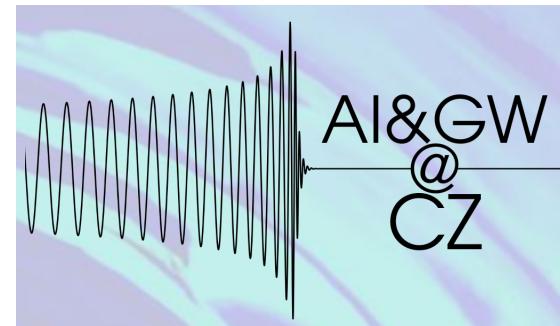


Modelling Extreme Mass Ratio Inspirals for LISA

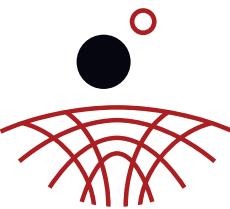
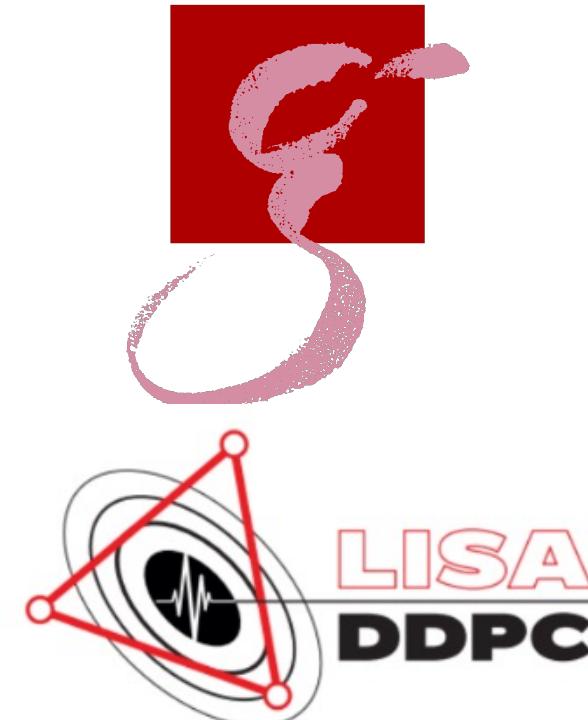
Overview, Challenges, and Roadmap

By Philip Lynch

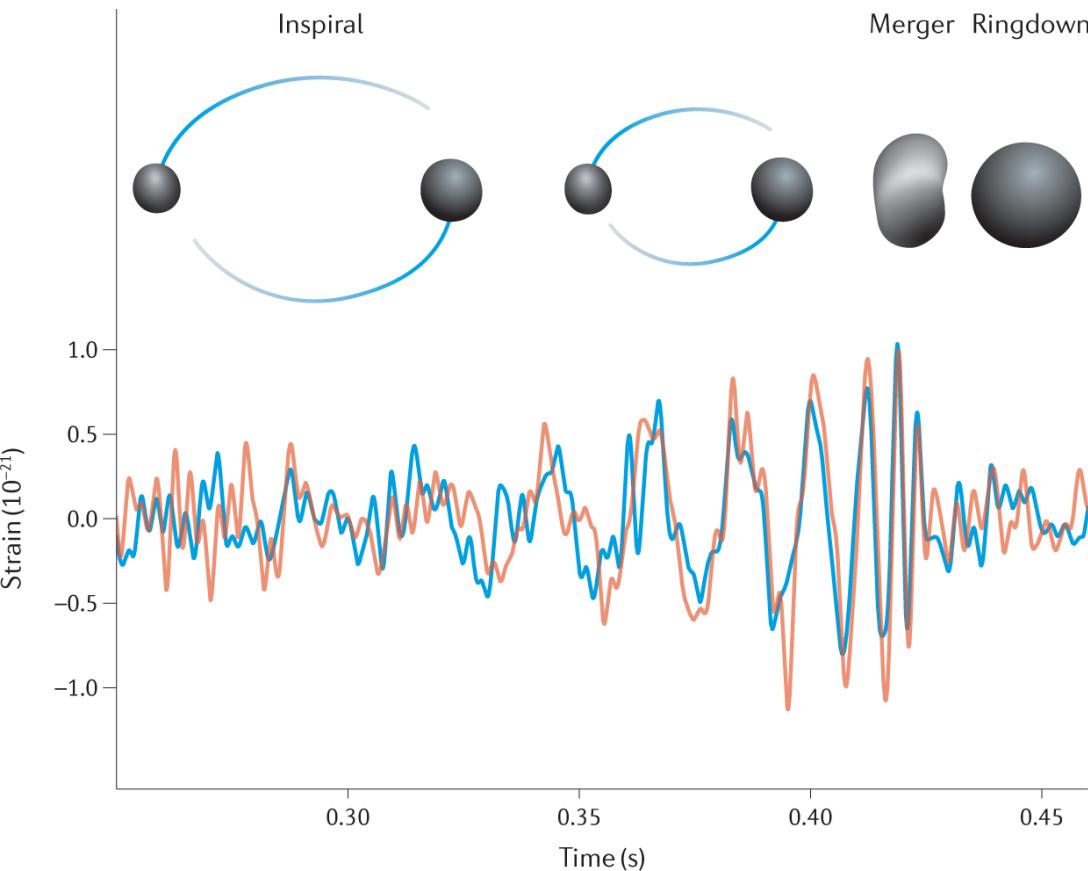


About me

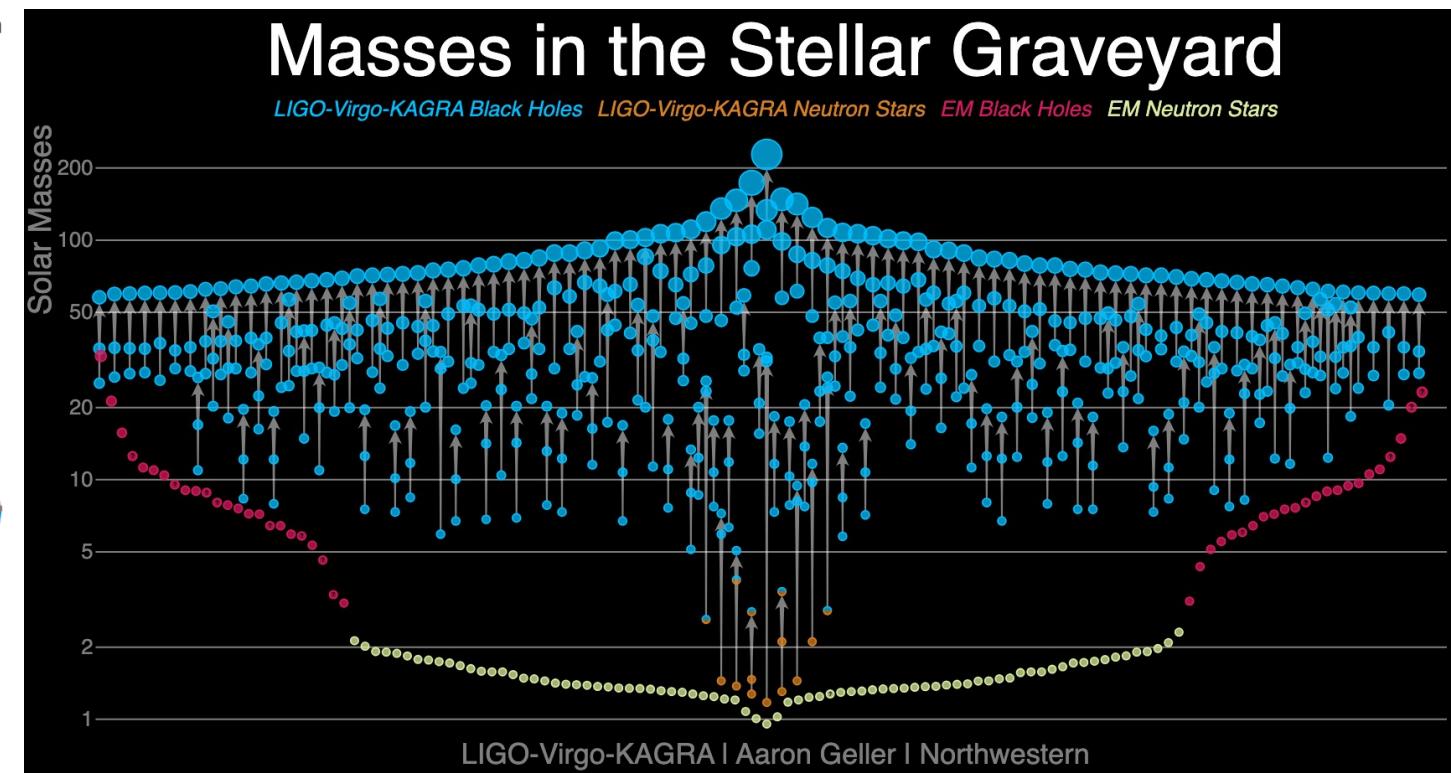
- 2014-18: **Undergrad in UCD**, Ireland
- 2018-22: **PhD in UCD** w/ Niels Warburton
(LISA Consortium Spokesperson)
- 2022- 2027: Postdoc in MPI for Gravitational Physics (AEI), Potsdam
- 2022: Joined **FastEMRIWaveforms dev team**
- 2024: Member of LISA's DDPC, CU Wav
- **Co-Lead of the EMRI Subunit** w/ Adam Pound

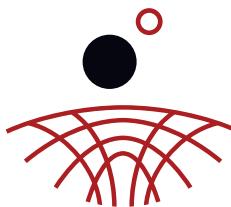


GW Astronomy



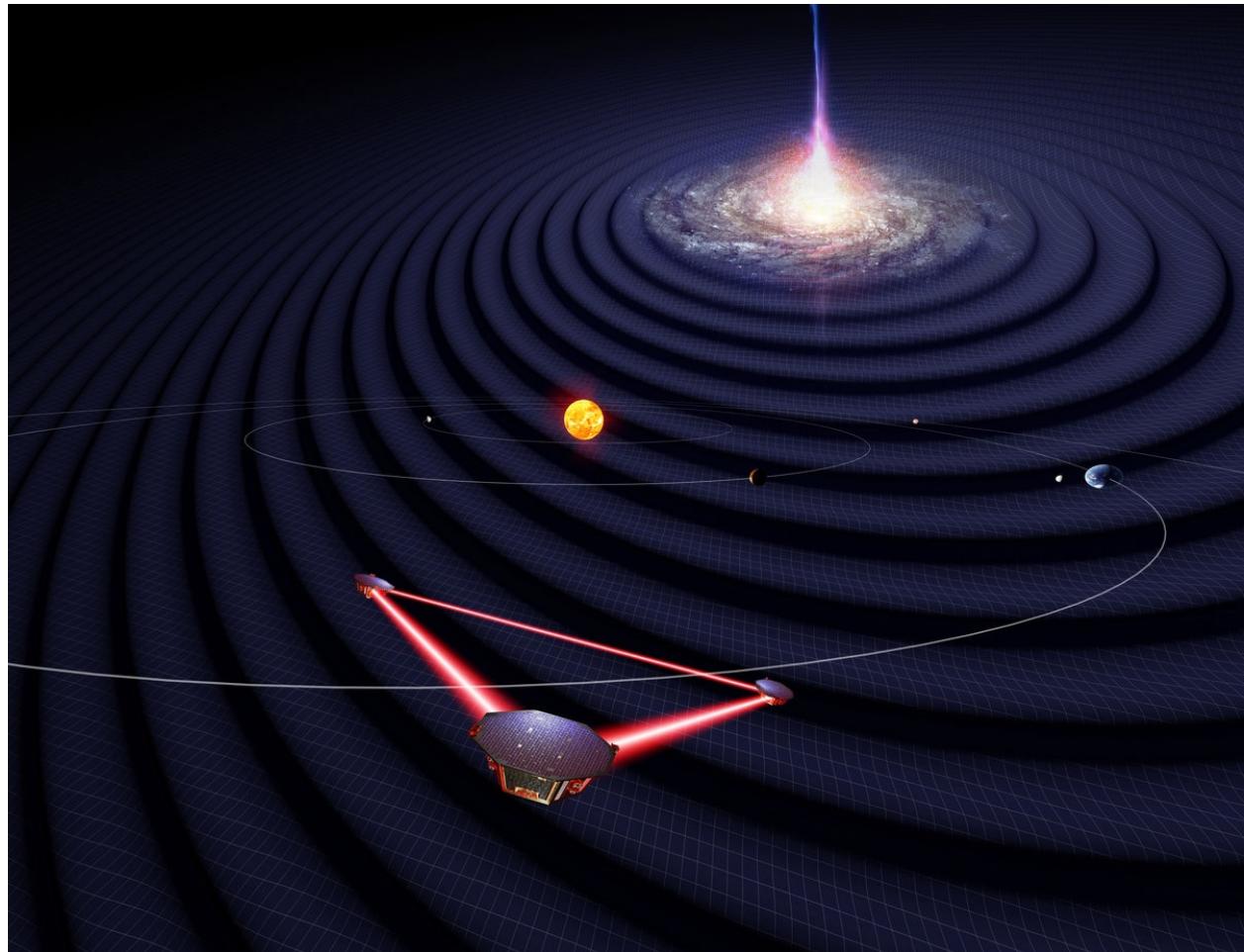
LVK Sensitivity: 10 – 10kHz





Space Based Detectors

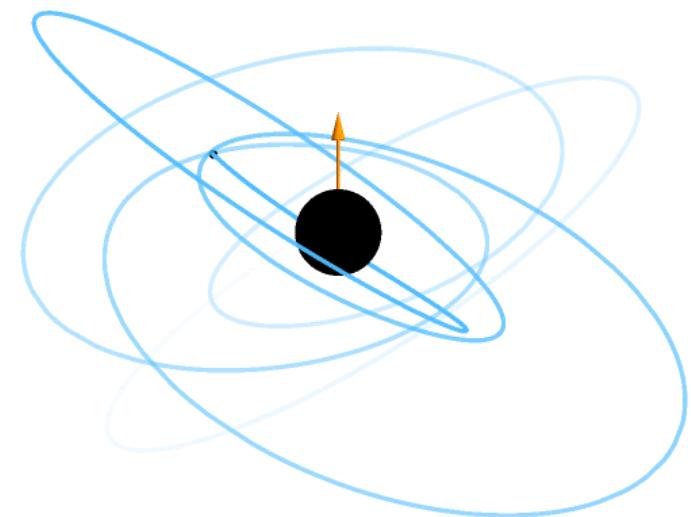
- LISA, TianQin & Taiji
- 2024: LISA formally **adopted by ESA**
- Scheduled to launch in **2035**
- Sensitive to **0.1 mHz to 1 Hz**
- See early inspiral of LVK sources
- **New GW Sources:** MBHBs, GBs, EMRIs & cosmological background etc.

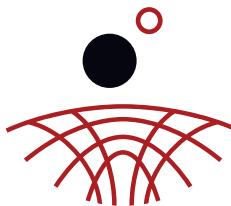




Extreme Mass Ratio Inspirals (EMRIs)

- **Massive black hole (MBH)** $m_1 \sim 10^5 M_\odot - 10^9 M_\odot$
- Compact object $m_2 \sim M_\odot - 10^2 M_\odot$
- **Small Mass Ratio**: $\epsilon = 10^{-4} - 10^{-7}$
- Loses energy and angular momentum to **GWs**
- Will stay in band for **years**!
- Expect ***eccentric, inclined, fast spinning*** primary $a \lesssim m_1$
- Expect **0–1000s** during LISA
- Accurate **MBH param est.** and stringent **tests of GR**



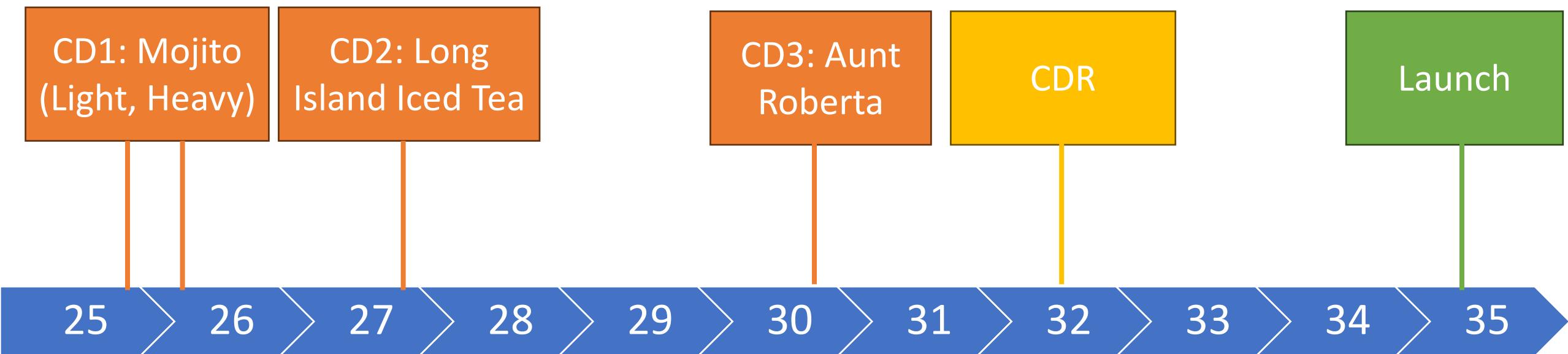


EMRI Waveform Model Requirements

- **Extensive**
 - Cover **full parameter space** and include all **necessary physics**
 - Primary spin, eccentricity, inclination, transient resonances & sec. spin
- **Fast**
 - Needed to do Bayesian Param est. MCMC -> Millions of WF evaluations
 - Generate a full WF in ideally $\sim 10ms$ but at least $< 1s$
- **Accurate**
 - Error in the phase < 1 radian over a 4 year inspiral
 - Can get away with larger errors in the amplitudes
 - Most of the SNR in the inspiral, don't need the merger ringdown
- **It has to be finished on time!**

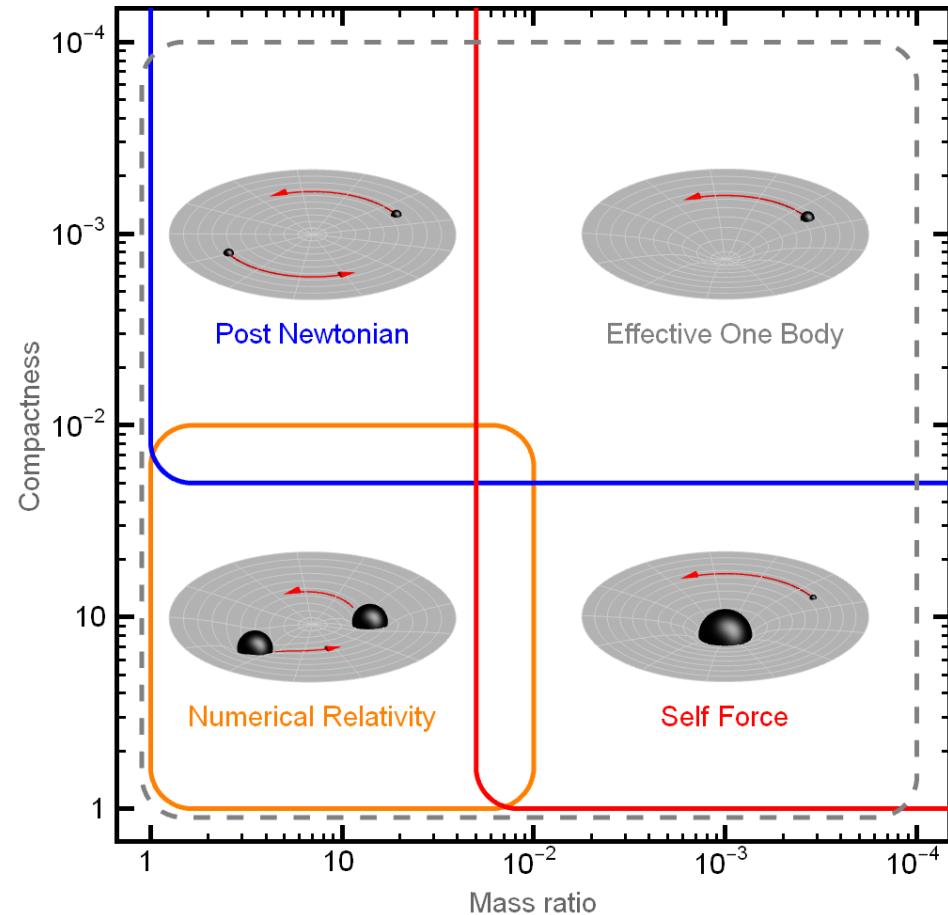


LISA Pre-Mission Timeline





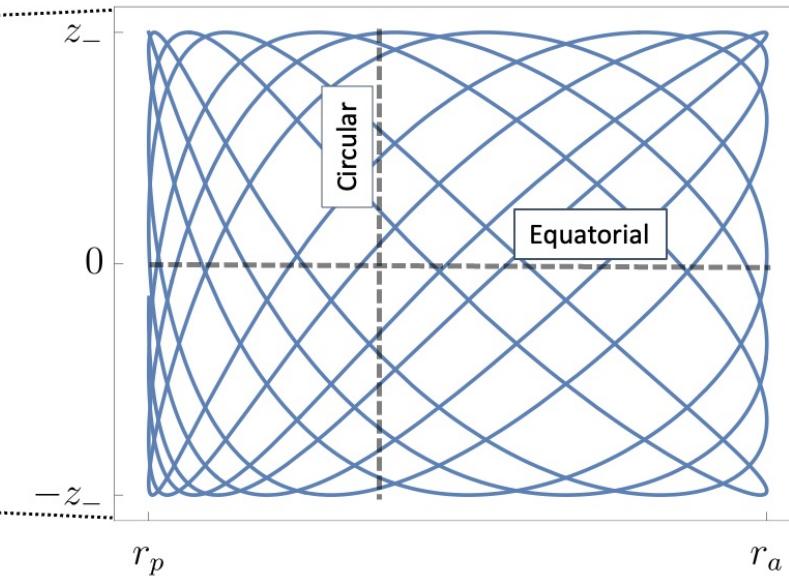
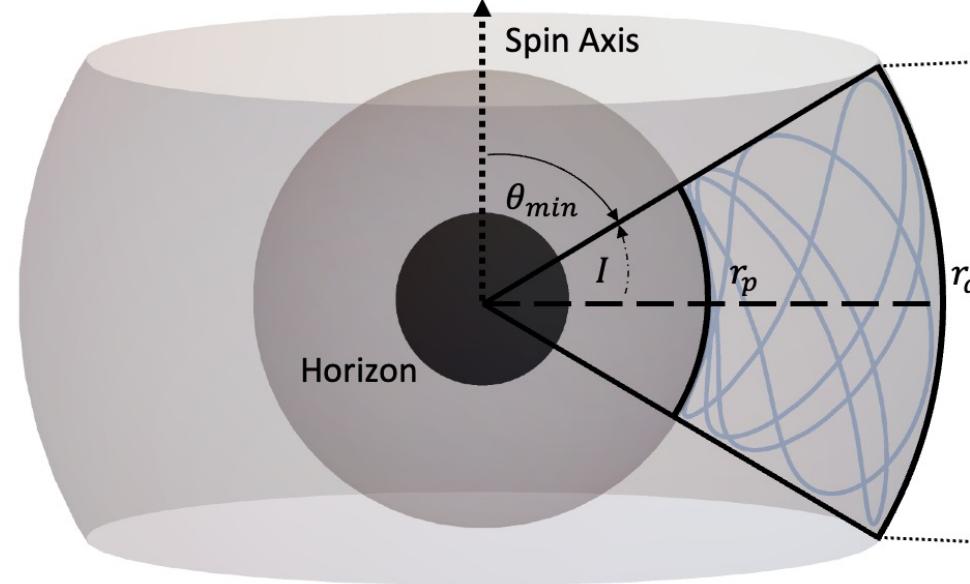
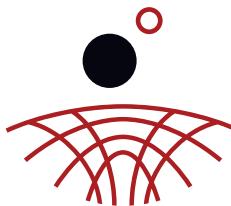
Relativistic Two Body Problem



- NR: Solve EFEs numerically
- Post-Newtonian: Expand in $(v/c)^2$
- Post-Minkoskian: Expand in G
- EOB: Use information from the above
- Self-force: Expand in $\epsilon = m_2/m_1 \in [0,1]$
- Symmetric mass ratio $\nu = \frac{m_2 m_1}{(m_1+m_2)^2} \in [0, \frac{1}{4}]$

Image Credit: M. van de Meent

Geodesic Motion in Kerr Spacetime

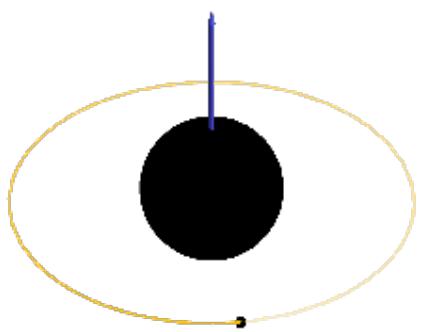


- **Orbital Elements P_j :** $p = \frac{2r_1 r_2}{M(r_1+r_2)}$, $e = \frac{r_1 - r_2}{r_1 + r_2}$, $x = \cos \theta_{inc} = \sqrt{1 - z_-^2}$
- **Orbital phases Φ_A :** $\frac{d\Phi_A}{dt} = \Omega_A^{(0)}(a, p, e, x)$
- Analytic solutions in terms of Jacobi Elliptic Integrals [Fujita & Hikida 09]
- Efficient conversions between Phases and Coordinates [PL & Burke 24]

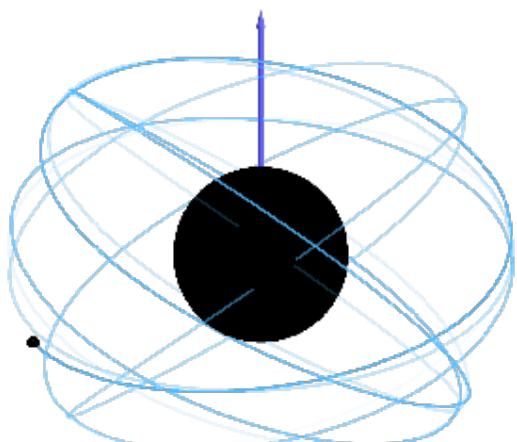
Equatorial
 $x = \pm 1$

Quasi-Circular

$$e = 0$$



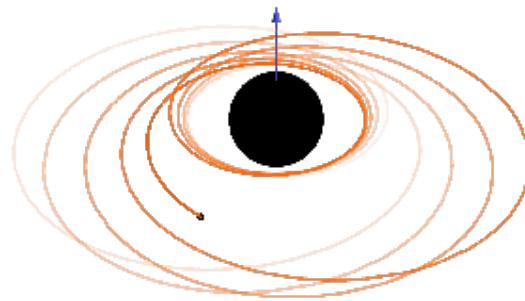
Inclined
 $-1 < x < 1$



(a.k.a Spherical)

Eccentric

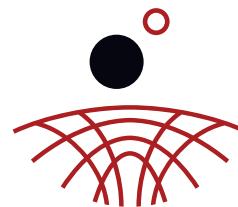
$$0 < e < 1$$



(a.k.a Generic)



Gravitational Self-Force

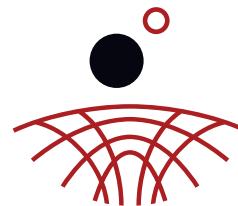


$$\frac{d^2 x^\alpha}{d\tau^2} + \Gamma_{\beta\gamma}^\alpha \frac{dx^\beta}{d\tau} \frac{dx^\gamma}{d\tau} = \epsilon a_{(1)}^\alpha + \epsilon^2 a_{(2)}^\alpha + \mathcal{O}(\epsilon^3)$$

Two-timescale analysis [Hinder & Flannigan 08]:

$$\varphi = \epsilon^{-1} \varphi^{(0\text{PA})} + \epsilon^{-1/2} \varphi^{(\text{res})} + \varphi^{(1\text{PA})} + \mathcal{O}(\epsilon)$$

- Adiabatic (0PA): $\langle a_{Diss}^{(1)} \rangle$ or \mathcal{E} & \mathcal{L} Fluxes
- Post-Adiabatic (1PA): $a^{(1)}$ & $\langle a_{Diss}^{(2)} \rangle$ & Sec. Spin
- Orbital Resonances: Generic Kerr only



Adiabatic Inspirals (OPA)

Teukolsky Equation

$$\mathcal{O} \psi_4 = 4 \pi \Sigma \mathcal{T}$$
$$\psi_4 = \frac{1}{2} \frac{d^2}{dt^2} (h_+ - i h_\times) \text{ as } r \rightarrow \infty$$

Solve in Frequency Domain with a geodesic source

Waveform Amplitudes

Multi-voice decomposition [Hughes +20]

$$h = \frac{m_2}{d_L} \sum_{lmkn} \mathcal{A}_{lmnk}(\vec{P})_{-2} Y_{lm}(\Theta, \Phi) \exp(-i \Phi_{mkn})$$

$$\text{where } \Phi_{mkn} = m \Phi_\phi + k \Phi_\theta + n \Phi_r$$

High dimensionality, lax accuracy requirement
ML techniques could be useful here

Asymptotic Fluxes

$$\langle j \rangle^\infty \text{ & } \langle j \rangle^{\mathcal{H}} \text{ where } \mathcal{J} = (\mathcal{E}, \mathcal{L}_z, \mathcal{Q})$$

Flux Balance Laws + Geodesic Relations

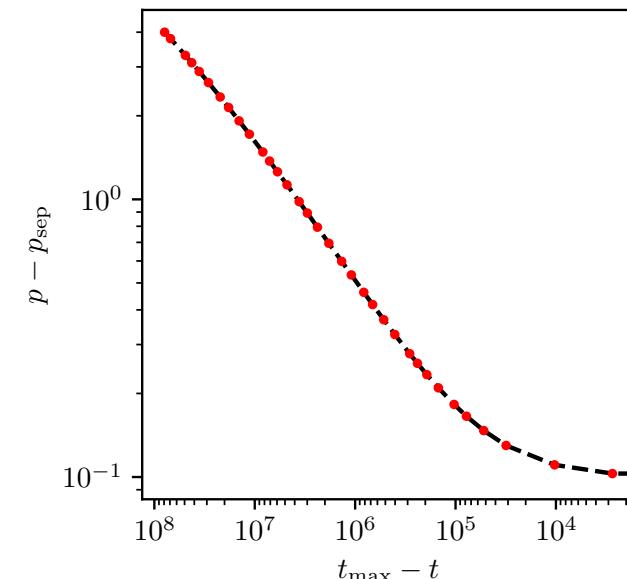
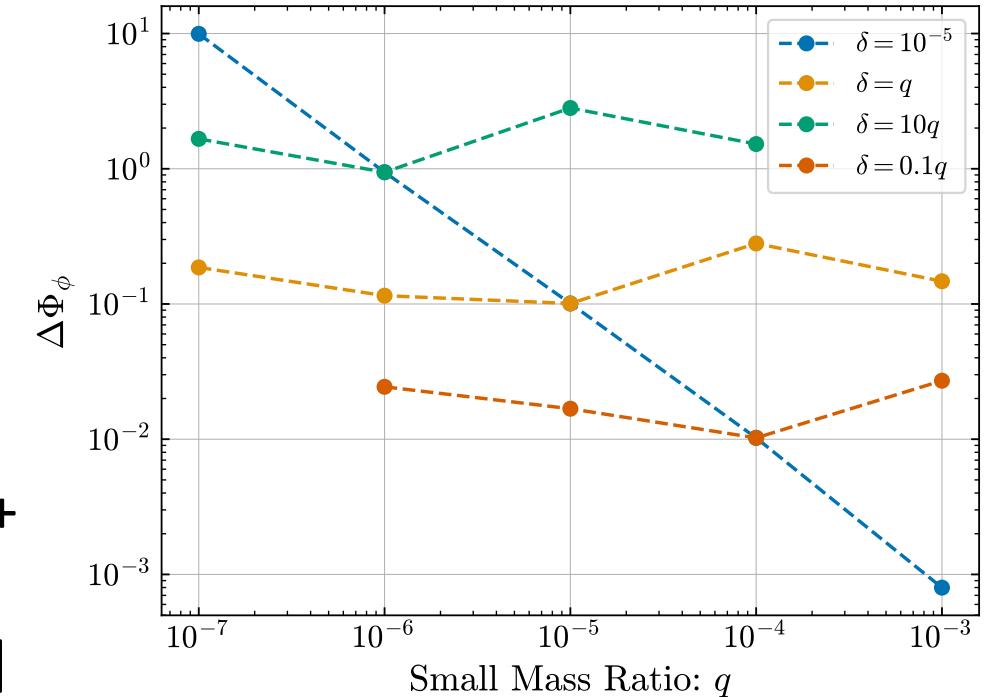
$$\rightarrow \dot{P}_j$$

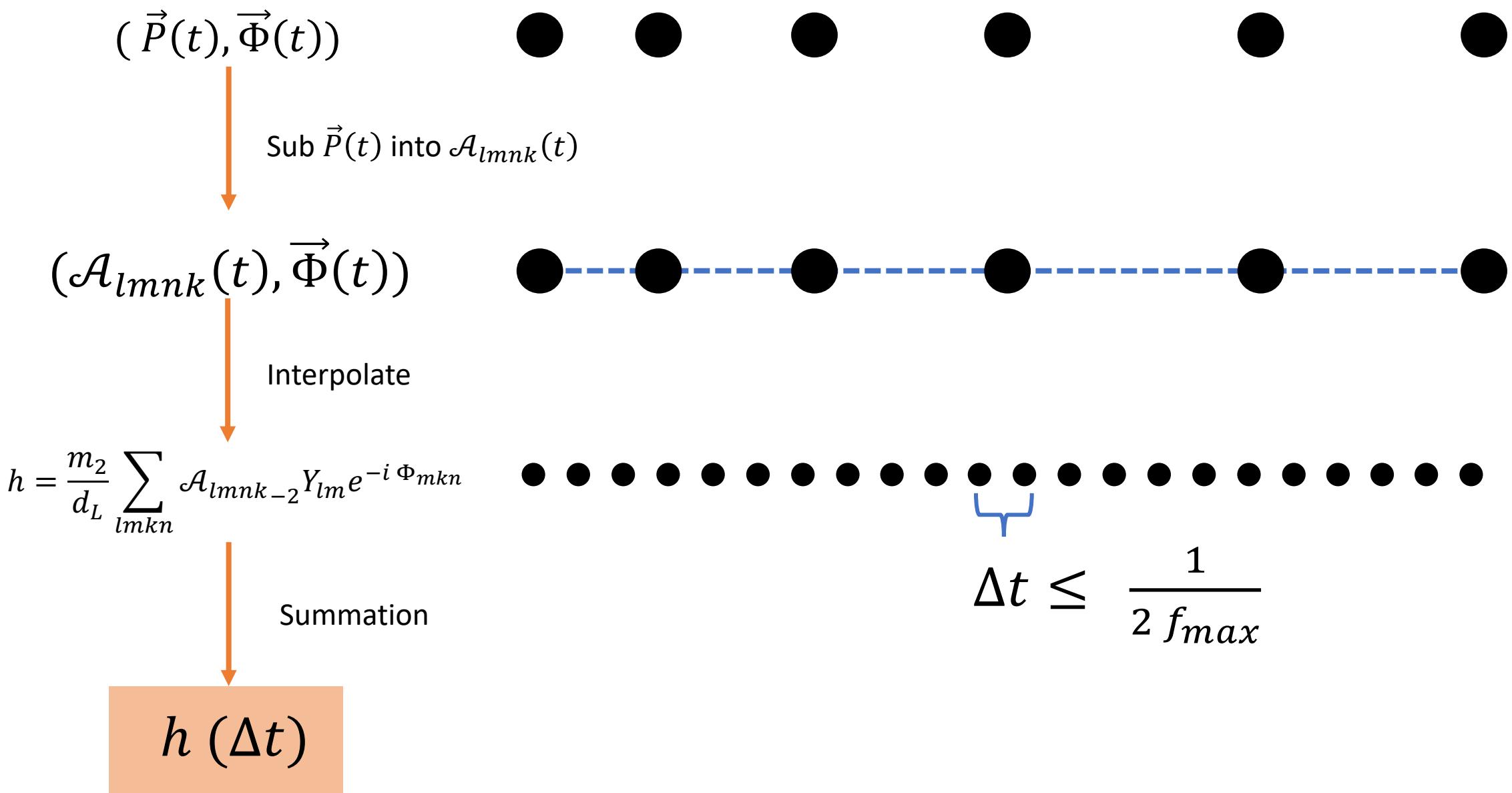
Low dimensionality, high accuracy requirement
ML will struggle

Interpolation of the fluxes

- Solve numerically at a single (a,p,e,x)
- **Offline step** (expensive, once):
 - Tile the parameter space and interpolate fluxes + mode amplitudes
 - **Flux rel. error** $\delta \leq \epsilon \leq 10^{-6}$ [Khavali, PL+ 2025]
 - Amplitudes can be much less accurate
- **Online step:** Numerically Solve ODEs

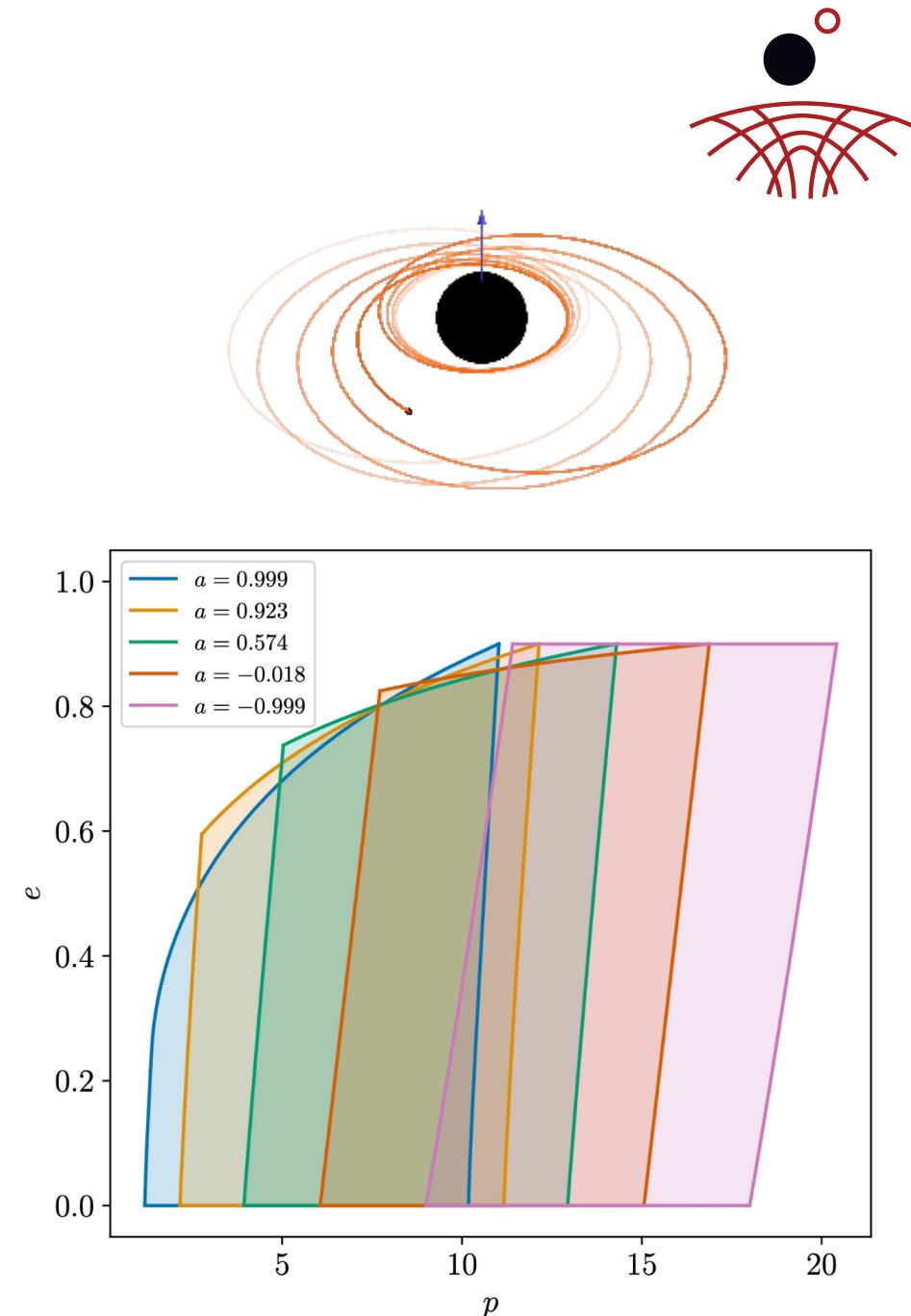
$$\begin{aligned}\dot{P}_j &= \nu F_j^{(1)}(\vec{P}) + \mathcal{O}(\nu^2) \\ \dot{\Phi}_A &= \Omega_A^{(0)}(\vec{P}) + \mathcal{O}(\nu)\end{aligned}$$





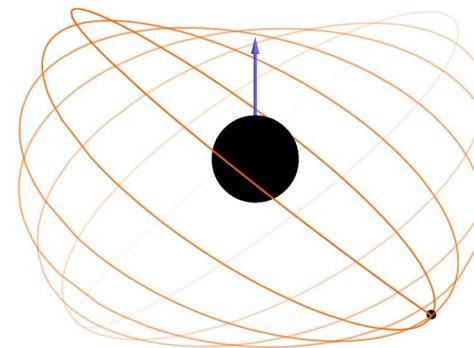
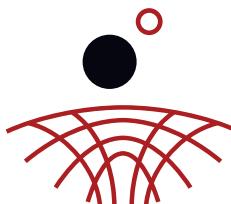
FastEMRIWaveforms (FEW)

- Leverages **GPUs** for Parallel calculations
- Can now do **eccentric** inspirals in **Kerr** [Chapman-Bird +(PL) 25]
- $a \in [-0.999, 0.999]$, $e \in [0, 0.9^*]$
- **Flux points:**
 - Inner (65 x 129 x 65)
 - Outer (33 x 65 x 33)
 - Total: 615,810 points costing 605,000 CPU hours
 - Interpolated with tri-cubic splines
- **Mode Amplitudes**
 - $\ell \in [2, 10]$, $m \in [0, \ell]$, $n \in [-55, 55]$
 - Total ~ 7000
 - Bi-cubic+linear spline
- **Mode Selector** only keeps necessary modes
- Speed: $\mathcal{O}(100ms)$

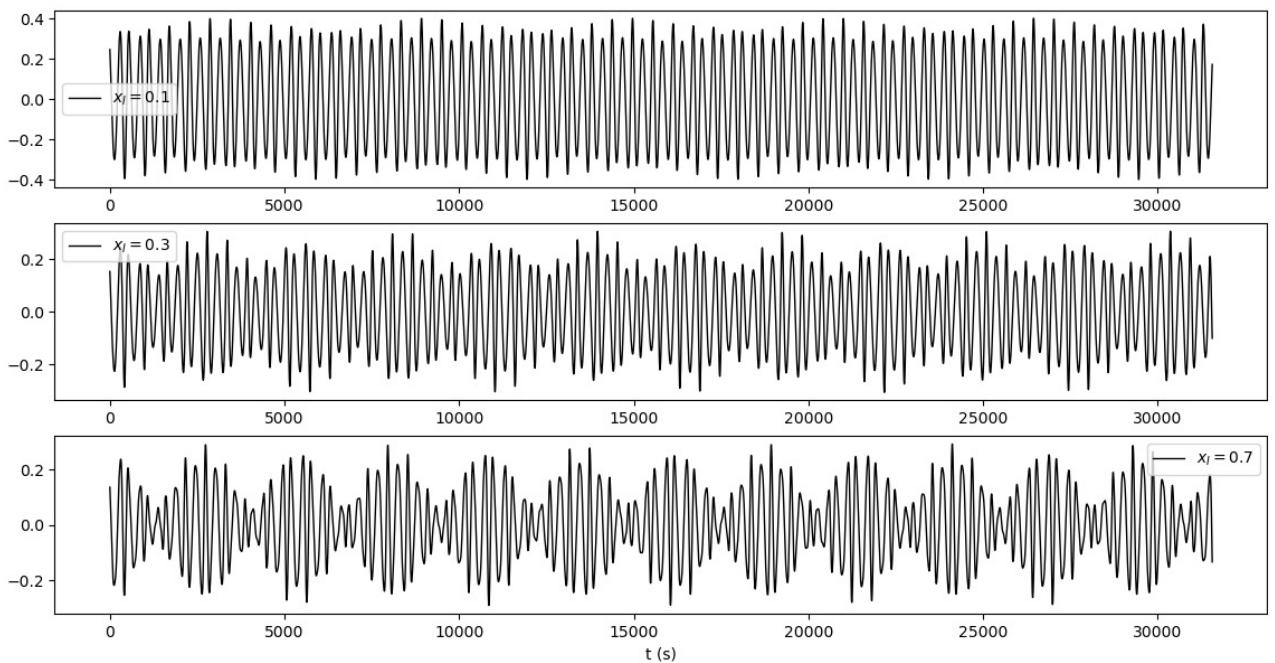


Spherical Orbits

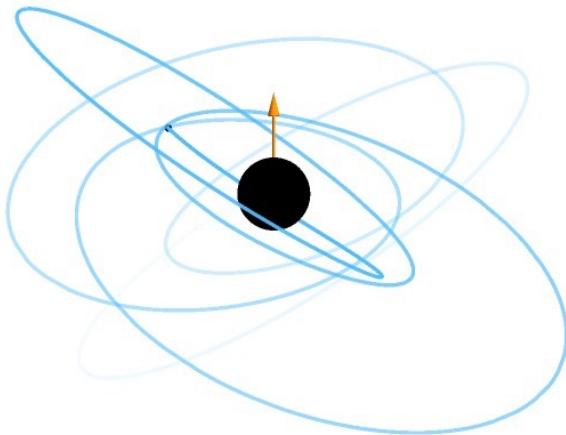
- FEW can now handle k modes
- Comparisons with high-order PN fluxes
- Cross checking interpolants with independent grids
- Checking spin precession conventions from comparable mass binaries (Leo Stein, Josh Mathews)
- Finish by Q1 2026



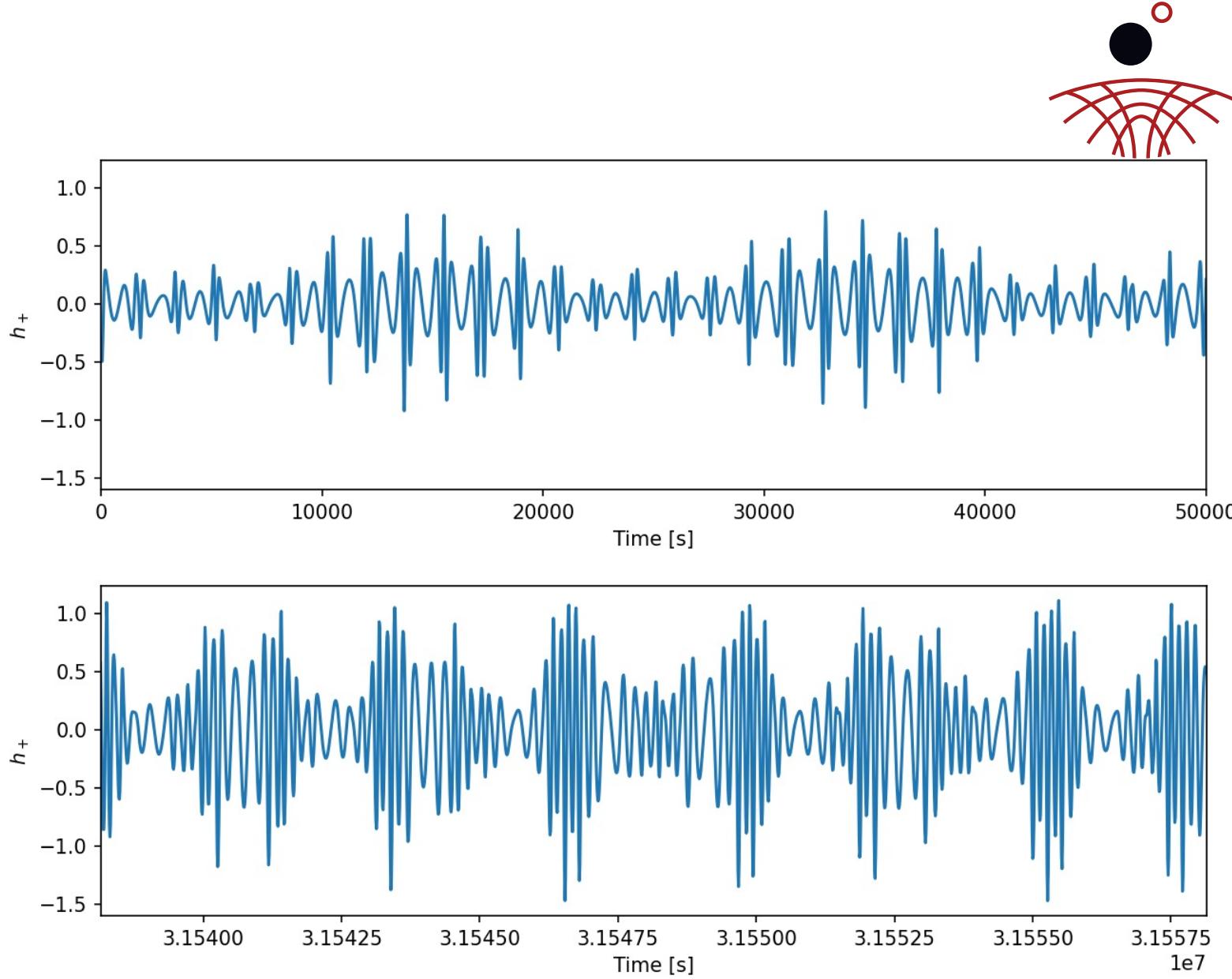
Varying initial inclination x_i (with $a=0.5$)

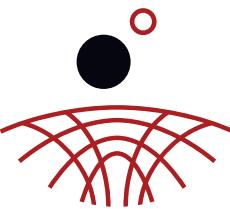


Generic Orbits

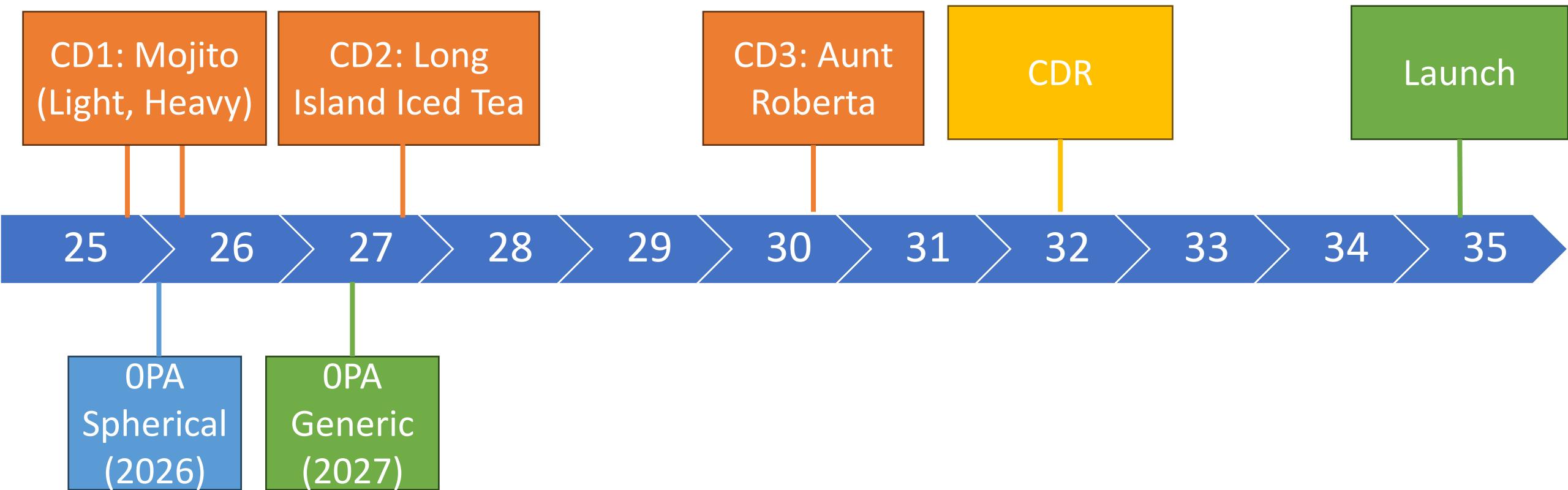


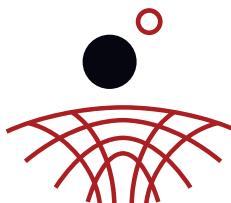
- Interpolated $a = 0.7$
- 4D spline interpolation now developed (Zach Nasipak)
- 4D Chebyshev interpolation developed (Philp Lynch)
- Still need to calculate the grid
- Aim for Q2 2027





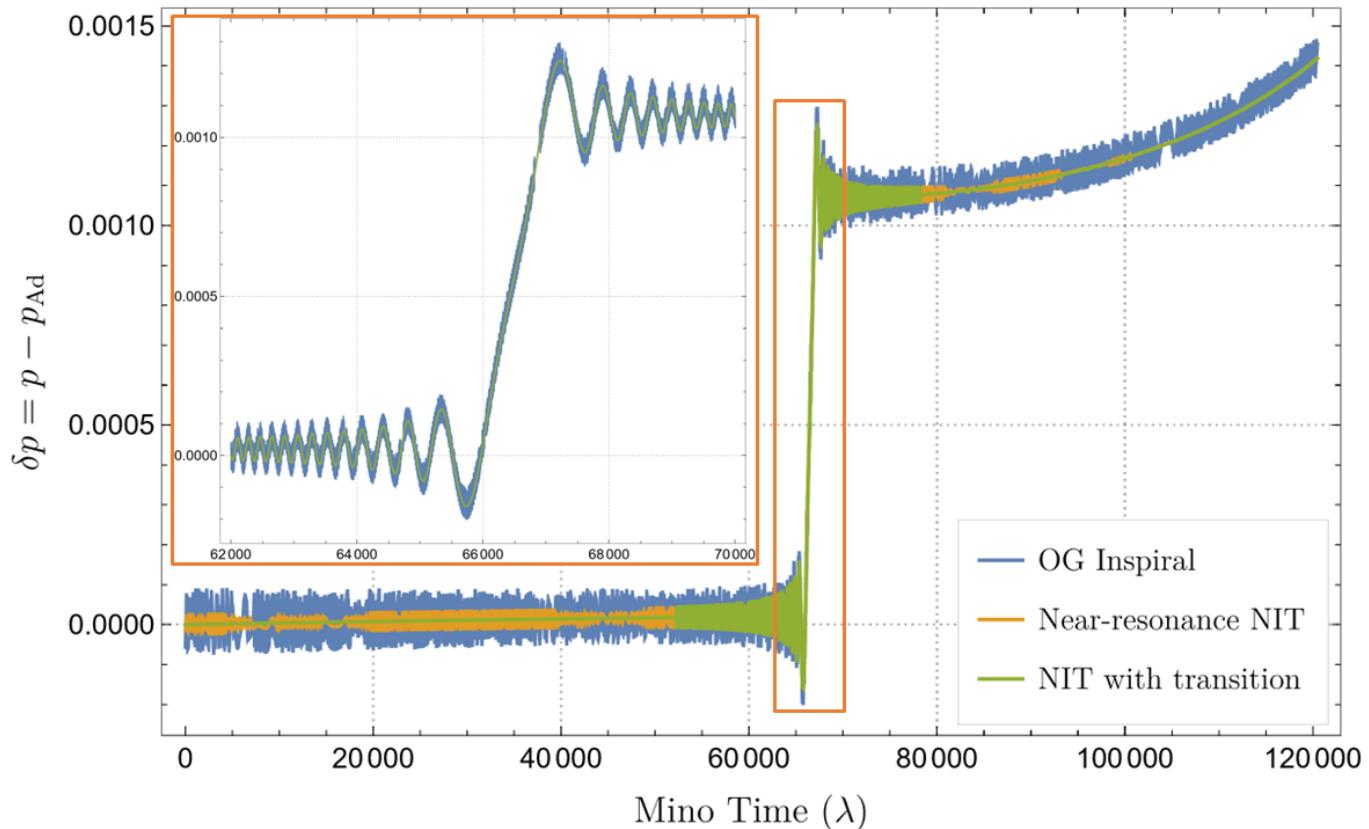
The FEWture Timeline





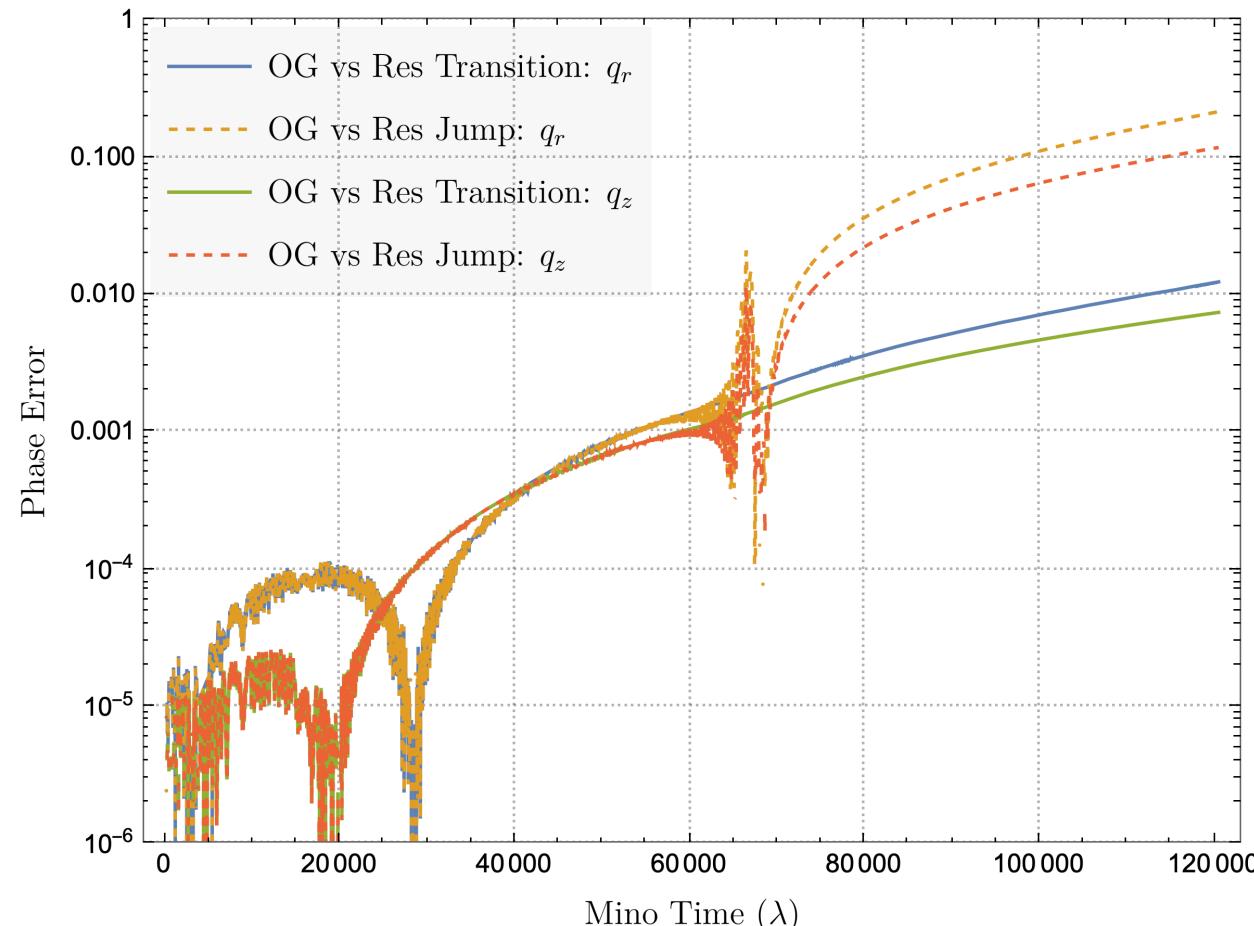
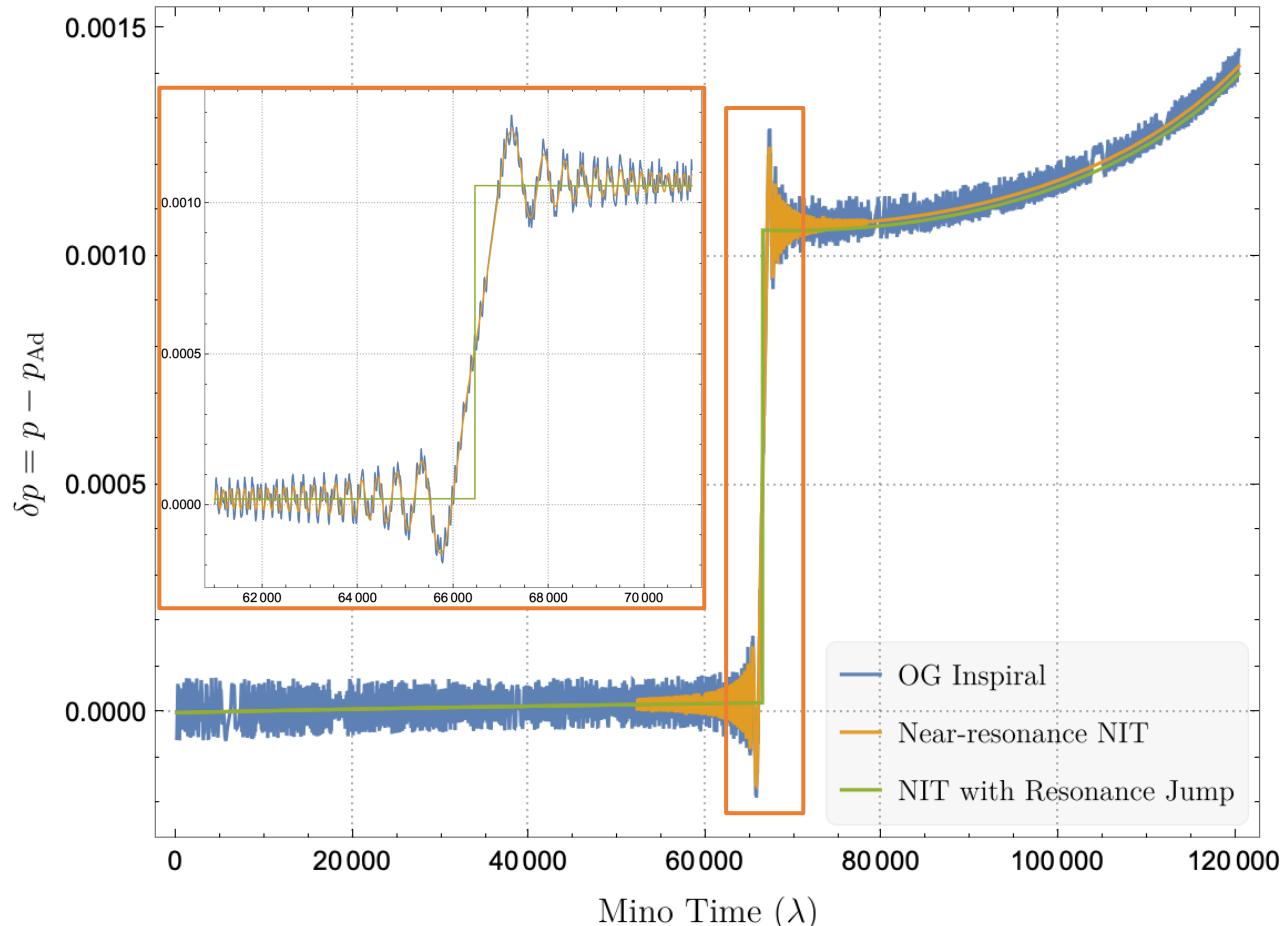
Transient Resonances ($\frac{1}{2}$ PA)

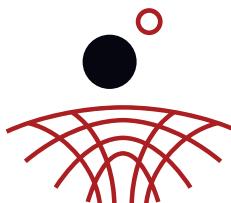
- Occur when $k \Omega_\theta - n \Omega_r = 0$ for $n, k \in \mathbb{N}$
- Cause “jumps” in $\vec{P} \propto \epsilon^{1/2}$
- Results in phase error $\propto \epsilon^{-1/2}$
- Work ongoing to model resonant jump in analytic Kludge model in FEW



Are Resonace Jumps Accurate enough for 1PA?

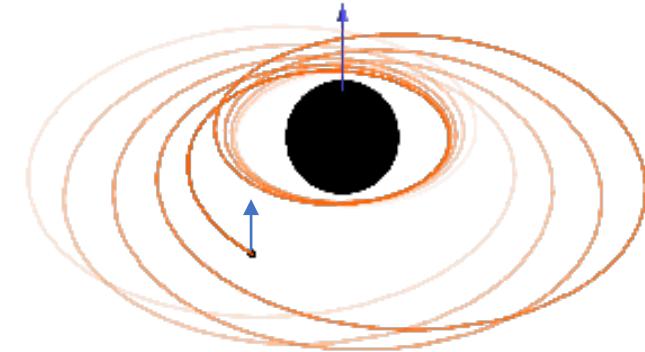
- Leading order Jump ΔP_j^{res} [Flannagan & Hinderer 12]
- Accurate Resonance treatment [Lynch+24]
- Vojtech and I are looking to find subleading piece needed for accurate 1PA inspirals



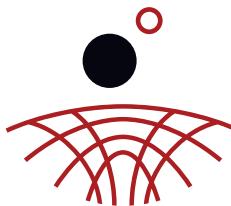


Spinning Secondary (1PA)

- For EMRIs, only need **linear in sec. spin**
- **Flux Balance Law** valid for spinning bodies in **generic orbits** [Grant 24]
- **Analytic Solutions** for Spinning Test Bodies [Skoupy & Witzany 25]
- Solve Teukolsky for a spinning test body and linearize in sec. spin
- Currently adding linear in spin corrections to **eccentric orbits** into FEW



$$\begin{aligned}\dot{P}_j &= \nu F_j^{(1)}(\vec{P}) + \nu^2 s_{\parallel} F_j^{(s)}(\vec{P}) \\ \dot{\Phi}_A &= \Omega_A^{(0)}(\vec{P}) + \nu s_{\parallel} \Omega_A^{(s)}(\vec{P})\end{aligned}$$



Second Order GSF (1PA)

- $g_{\alpha\beta}^{exact} = g_{\alpha\beta}^{(0)} + \epsilon h_{\alpha\beta}^{(1)} + \epsilon^2 h_{\alpha\beta}^{(2)}$
- Split into Regular and Singular pieces
- Sub into EFEs: $G_{\alpha\beta}[g] = 8\pi T_{\alpha\beta}$

- ϵ^0 : $G_{\alpha\beta} \left[g_{\alpha\beta}^{(0)} \right] = 0$

- ϵ^1 : $G_{\alpha\beta}^{(1)} \left[h_{\alpha\beta}^{(1)R} \right] = T_{\alpha\beta}^{(1)}[z] - G_{\alpha\beta}^{(1)} \left[h_{\alpha\beta}^{(1)S} \right]$

- ϵ^2 : $G_{\alpha\beta}^{(1)} \left[h_{\alpha\beta}^{(2)R} \right] = T_{\alpha\beta}^{(2)}[z] - G_{\alpha\beta}^{(2)} \left[h_{\alpha\beta}^{(1)}, h_{\alpha\beta}^{(1)} \right] - G_{\alpha\beta}^{(1)} \left[h_{\alpha\beta}^{(2)S} \right] - \partial_{\tilde{t}} h_{\alpha\beta}^{(1)}$

$$\dot{P}_j = \nu F_j^{(1)}(\vec{P}) + \nu^2 F_j^{(2)}(\vec{P})$$

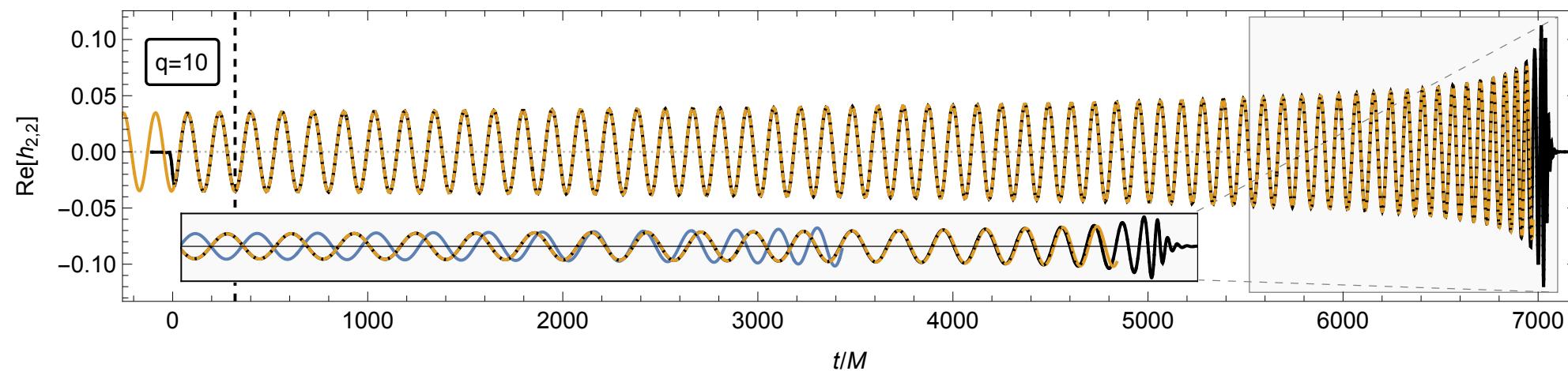
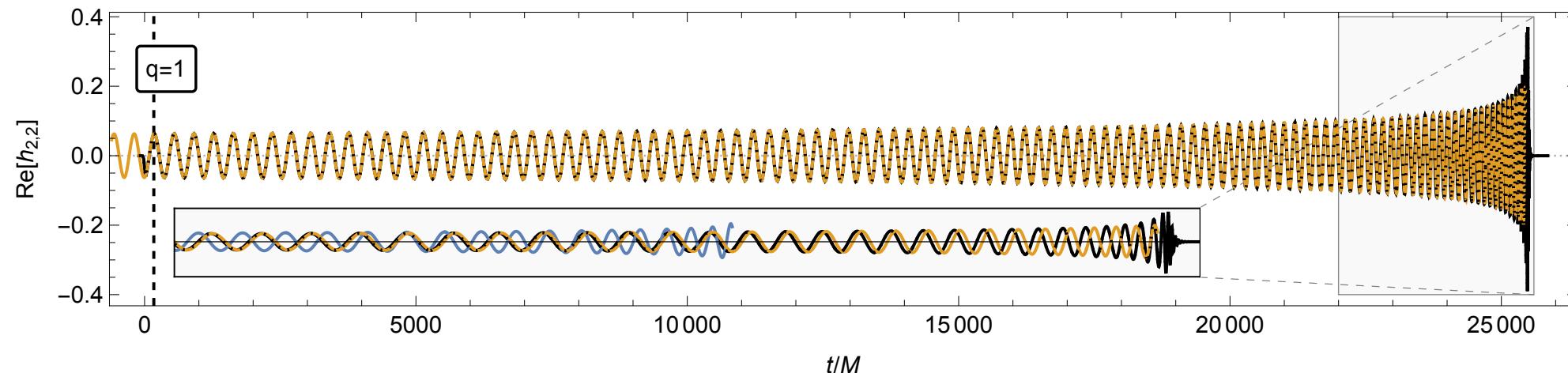
$$\dot{\Phi}_A = \Omega_A^{(0)}(\vec{P}) + \nu \Omega_A^{(1)}(\vec{P})$$

This is very hard – Barry Wardell (2025)

- Multiscale Self-Force Collaboration

Quasi-circular non-spinning

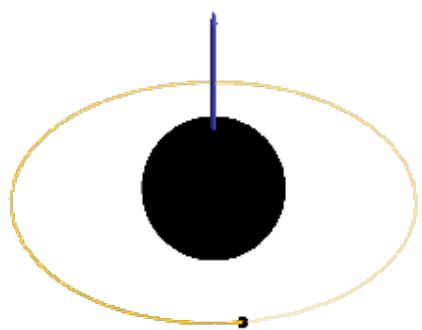
[Warburton & Pound & Wardell + 19,20,21]



Equatorial
 $x = \pm 1$

Quasi-Circular

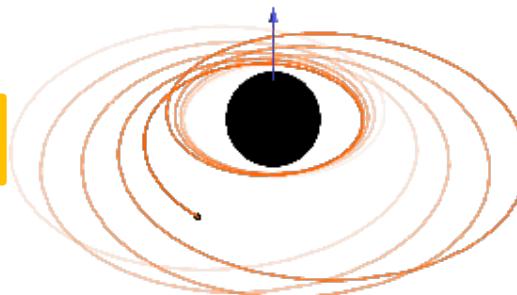
$$e = 0$$



1PA: Q2 2027

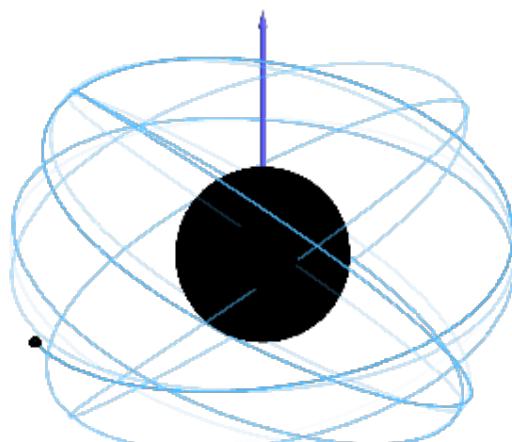
Eccentric

$$0 < e < 1$$



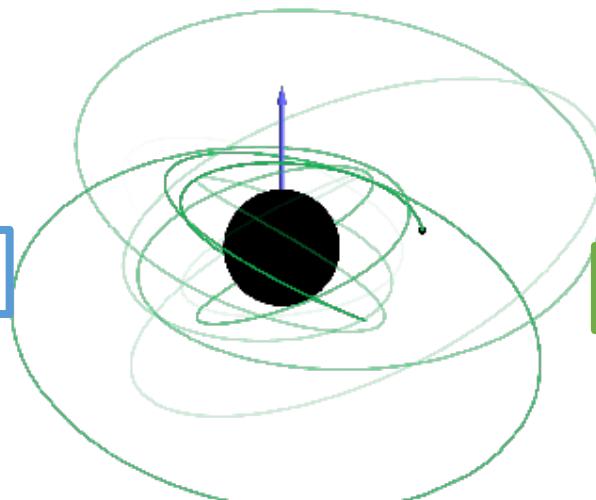
1PA ~ 2029 ?

Inclined
 $-1 < x < 1$



1PA: Q2 2028

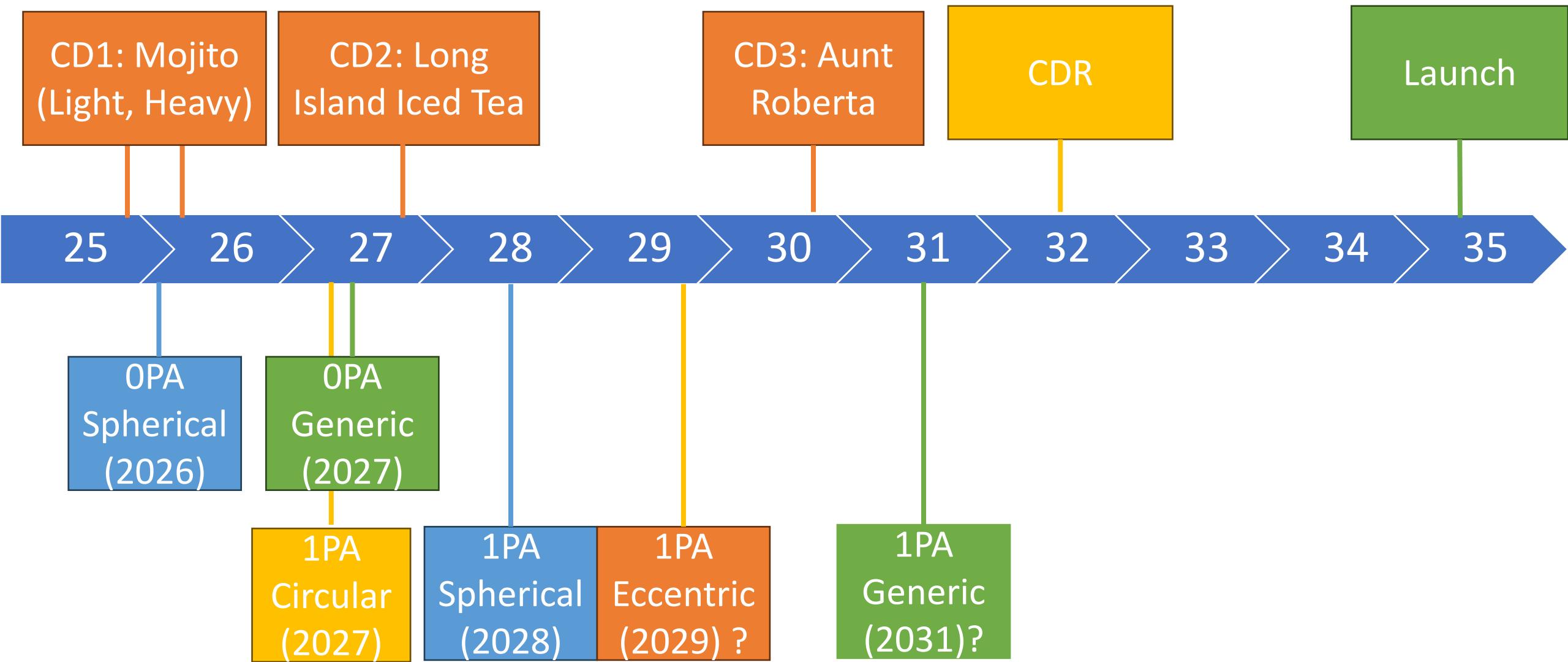
(a.k.a Spherical)



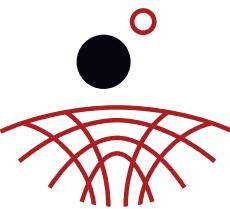
1PA ~ 2030??

(a.k.a Generic)

The FEWture Timeline



1PA Quasi-circular Schwarzschild in FEW [Burke+ (PL) 24]



$$\dot{P}_j = \nu F_j^{(1)}(\vec{P}) + \boxed{\nu^2 F_j^{(2)}(\vec{P})} + \nu^2 s_{\parallel} F_j^{(s)}(\vec{P})$$

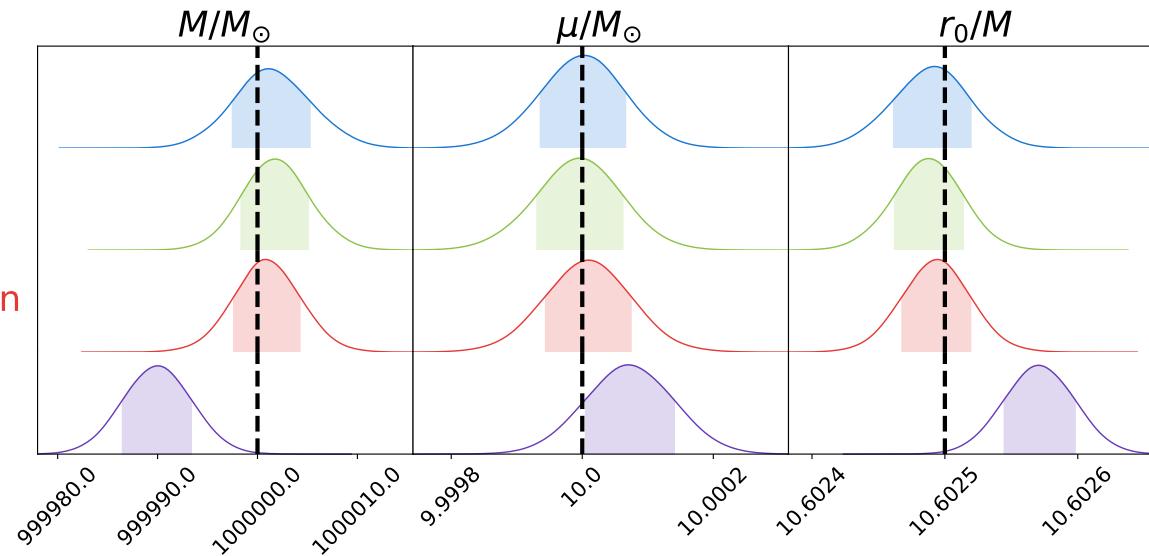
$$\dot{\Phi}_A = \Omega_A^{(0)}(\vec{P}) + \nu \Omega_A^{(1)}(\vec{P}) + \nu s_{\parallel} \Omega_A^{(s)}(\vec{P})$$

cir1PA w/ spin

cir1PA w/o spin

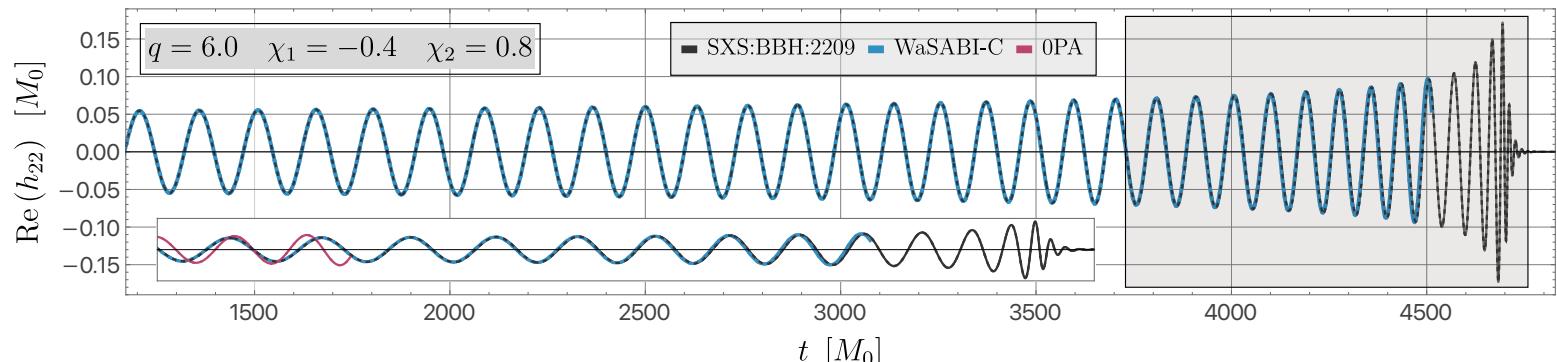
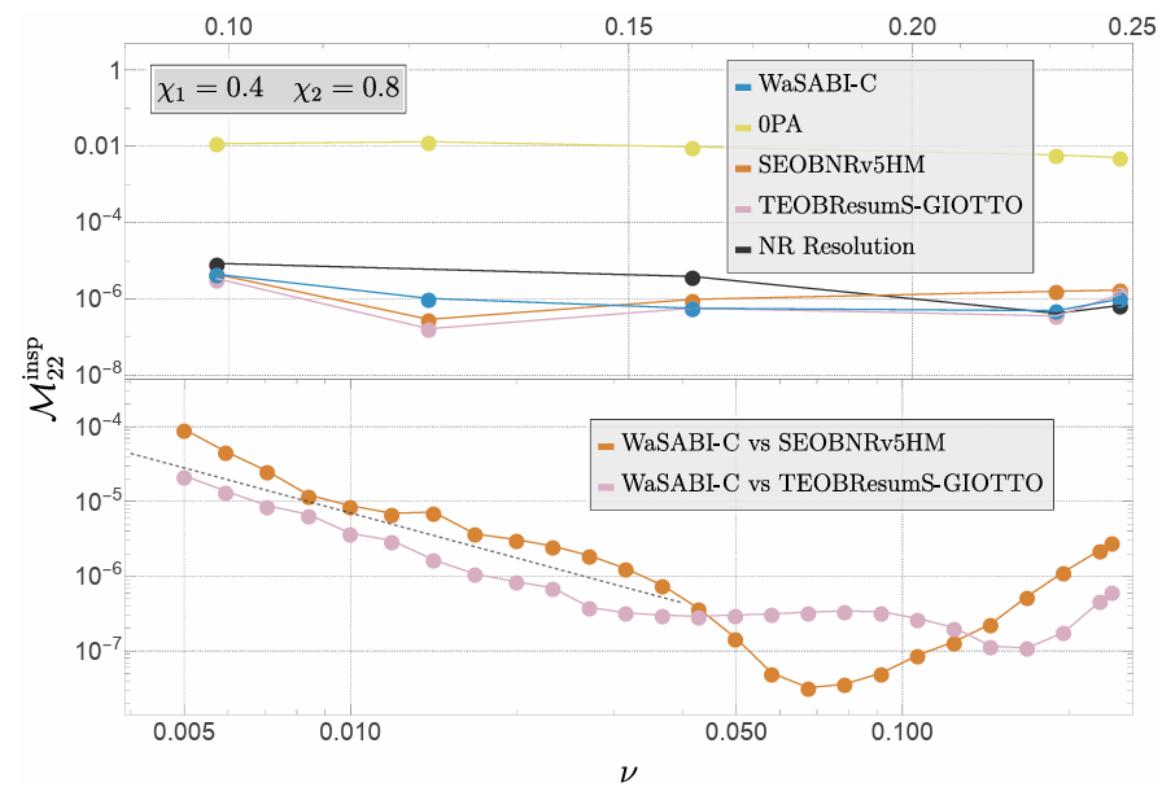
cir0PA + 1PA-3PN w/o spin

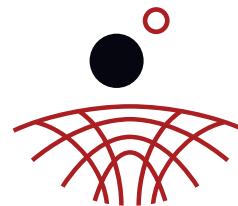
cir0PA w/o spin



Hybrid waveforms w/ PN

- Failsafe if we do not get complete 1PA
- First published model:
 - Slow spinning primary
 - Aligned Secondary Spin
 - 1PA and higher terms extracted from PN
- In EMRI limit: about 0.1 rad dephasing relative to complete 1PA model in nonspinning case
- Mathematica implementation: WaSABI





Summary

Adiabatic OPA:

- FEW: $\mathcal{O}(100ms)$
- On track to cover spherical & generic by 2027
- Can easily add 1PA results

Resonace $\frac{1}{2}$ PA:

- Implementation in Kludge
 - May need to go past leading order jumps

Spinning Secondary 1 PA:

- Theory is done
- Starting on the practical
- Can be added modularly

2nd Order SF (1 PA):

- Slow but steady progress
- Spherical before eccentric
 - May need PN hybrids